### Extended MHD Simulations of D3D 86144<sup>1</sup>

T.A. Gianakon

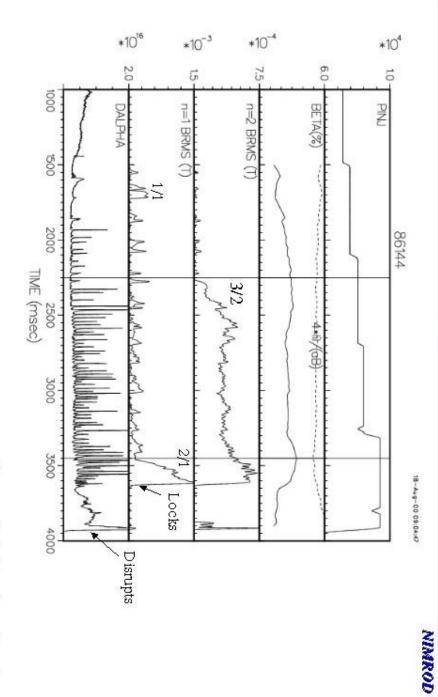
D.D. Schnack

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Long Beach, CA

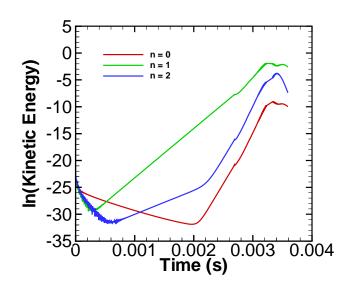
<sup>&</sup>lt;sup>1</sup>Research supported by U. S. DoE

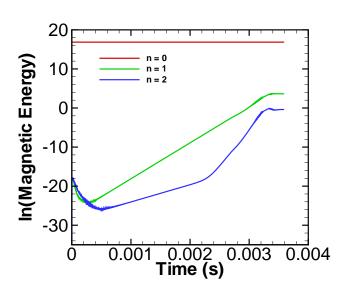
### Modern Tokamak Discharges Have Rich Magnetic Behavior



Courtesy of R. LaHaye, General Atomics

### Both n=1 and n=2 are unstable. Mode coupling dominates n=2.

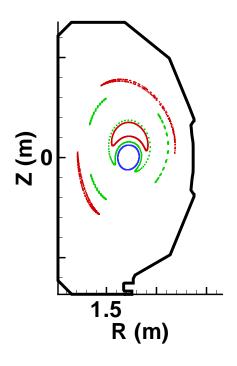


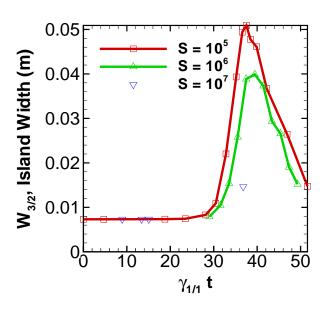


- Change in n=2 slope is nonlinear coupling 1/1 to 2/2.
- NIMROD growth rate is  $\gamma = 4.58(10^3)(s^{-1})$  at  $S = 10^7$  and  $P_m = 1000$ .
- Experimental measurement from a fit of the Mirnov coils gives  $\gamma = 1.68(10^4)(s^{-1})$ .

### Secondary islands form, but too small to explain 3/2.

- Secondary island formed by nonlinear couplings.
- Secondary island width is reduced with decrease in resistivity (S increases).





## Extended MHD seeks numerically tractable forms to close MHD equations.

• Resistive MHD is incomplete.

is too small. A 3/2 tearing instability is observed in the experiment, but with resistive MHD the island width

- ullet The |B| variations along a magnetic field lead to particle trapping effects.
- Particle trapping cause neoclassical effects:

Damping of poloidal flows.

Neoclassical enhancement of the polarization current.

Generation of bootstrap currents.

• A simplified viscous stress tensor captures most neoclassical effects

$$abla \cdot ec{\Pi}_lpha = m_lpha n_lpha \mu_lpha \langle B_0^2 
angle rac{ec{v}_lpha \cdot 
abla heta}{(ec{B}_0 \cdot 
abla heta)^2} 
abla heta,$$

where  $\alpha$  indicates electrons and ions, and  $\mu_{\alpha}$  is a flow damping rate.

• Extensive testing of more complete closures completed to verify the simple

<sup>&</sup>lt;sup>2</sup>T.A. Gianakon, S.E. Kruger, C.C. Hegna, Physics of Plasmas, **9** January 2002.

# The perturbed bootstrap current can cause instability.

• Pressure flattening causes a perturbed bootstrap current.

Anisotropic heat diffusion about an island causes  $\nabla P \longrightarrow 0$ .

A helical hole in the bootstrap current is formed.

The hole in the perturbed bootstrap current causes island growth.

Physics captured by a simple Rutherford island evolution equation.

$$rac{k_0}{\eta^*}rac{dW}{dt} = \Delta^{'} + rac{W_{nc}}{W}rac{W^2}{W^2 + W_d^2}$$

The  $\Delta'$  is nominally negative, (standard tearing is positive.)

The first half of the second term is the bootstrap current.

The second half is the correction for anisotropic heat diffusion.

ullet The anisotropic heat diffusion causes a nonlinear threshold for instability.

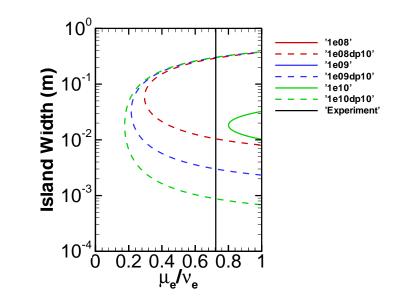
For  $W < W_d$ , no pressure flattening occurs. For  $W > W_d$ , pressure flattening occurs.

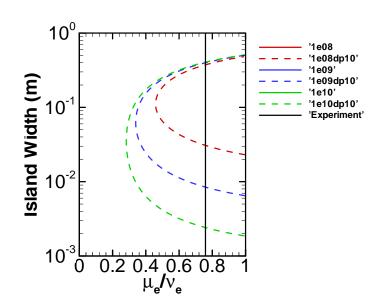
The formation of secondary islands is necessary to seed this instability, the neoclassical tearing mode (NTM).

### Seed islands cross NTM stability boundaries for 86144.

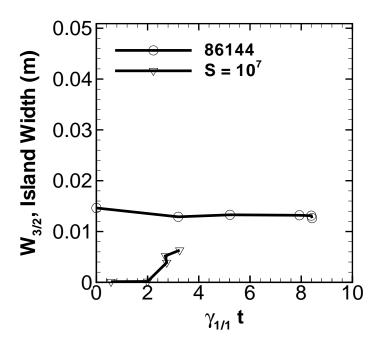
The 3/2 stability boundaries.

The 2/1 stability boundaries.





- ullet Uncertainty in both  $\chi_{||}/\chi_{\perp}$  and  $\Delta^{'}$  exist.
- ullet The 2/1 is not neoclassically unstable—consistent with experiment.
- The 3/2 is potentially neoclassically unstable.



### Summary

- Nonlinear simulations with NIMROD of DIIID Shot 86144.02250 illustrate a rich array of MHD and extended MHD effects.
- Growth and saturation of an internal kink can be identified.
- Secondary island formation from the internal kink is modeled.
- Growth of neoclassical tearing modes.